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Dynamic Modeling and Analysis of a Novel 3-RRR Parallel Shoulder

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ABSTRACT

The dynamic analysis of a novel 3-DOF robot shoulder joint based on orthogonal spherical 3-RRR parallel mechanism was presented. The kinematics constraint equation of this mechanism was established based on the geometric structure and the kinematics model of moving platform was derived. The dynamics model was established based on Lagrange method and the effective inertia, coupling inertia and driving torque of the shoulder joint were analyzed. The variation of the dynamics parameters with the change of the mechanism was discussed and the dynamic coupling relationship between branches was analyzed. The analysis results show that the posture change of the mechanism has a great influence on the dynamics parameters in the process of movement. The results will be useful to improve the control scheme of this mechanism and the selection of servo motor.

Key words: Shoulder, dynamic modeling, virtual work principle

INTRODUCTION

Spherical 3-DOF (degrees of freedom) parallel mechanism is a limited-DOF parallel mechanism which can be widely used for different kinds of practical applications (Zhen *et al.*, 2006). It both maintains the inherent advantages of parallel mechanisms and possesses several other advantages in terms of simple structure, low cost in manufacturing and operations. Therefore it has a potentially wide range of applications, such as motion simulator, bionic joint and machine tools. It is attracting the attention of the various researchers and many prototypes based on spherical 3-DOF parallel manipulator have been developed, such as agile eye (Gosselin and Hamel, 1994; Gosselin and St-Pierre, 1997), NC rotary table (Zeng *et al.*, 2001), wrist joint (Sun *et al.*, 2003), shoulder joint (Jin and Rong, 2007), waist joint (Li and Jin, 2007), bionic eye (Sellaouti *et al.*, 2002) etc. The dynamics is an important research content of robot which plays an important role in the, trajectory planning, driver selection and control scheme of the robot (Craig, 2011) and the establishment of the dynamic model for mechanism is the premise and basis for the dynamics analysis of parallel mechanism (Feng *et al.*, 2006; Ji *et al.*, 2012; Chen *et al.*, 2013).

On the basis of the optimal scale parameter values, a novel 3-DOF robot shoulder joint based on orthogonal spherical 3-RRR parallel mechanism is designed but the dynamics of the novel 3-DOF robot shoulder joint is not discussed (Zhang *et al.*, 2015), so the dynamic analysis of the robot shoulder joint is presented in this study. Through the analysis and calculation of kinetic energy and potential energy of each component, the dynamics model of the relationship between

driving torque and the motion of the system parameters is established based on Lagrange method. On the basis of the dynamics model, the effective inertia, coupling inertia and driving torque of the parallel mechanism system are analyzed.

MATERIALS AND METHODS

Description of the shoulder: The computer aided design model and schematic of the shoulder joint is shown in Fig. 1 and 2, respectively. The prototype of the shoulder joint in the study is a novel parallel 3RRR orthogonal spherical parallel mechanism. It consists of a fixed platform, a moving platform and three limbs which consist of a framed link and a connecting rod with identical structure. Each limb that consists of three revolute joints in series connects the fixed platform to the moving platform. Thus, the moving platform is attached to the fixed platform by three identical

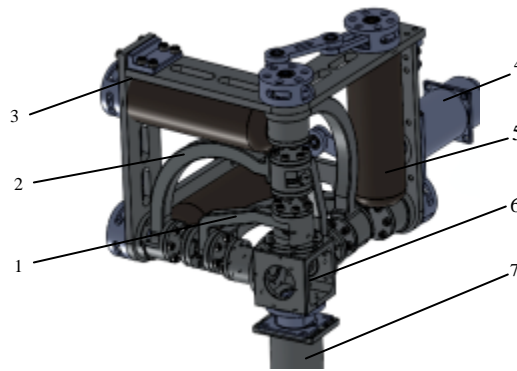


Fig. 1: Prototype of shoulder joint, 1: Connecting rod, 2: Frame link, 3: Frame, 4: Trunk connector, 5: Actuator, 6: Moving platform, 7: Big arm

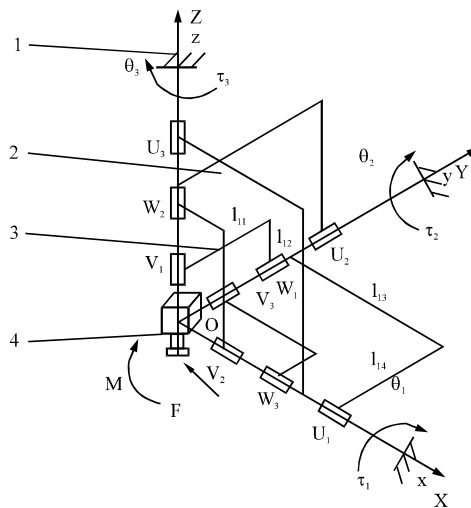


Fig. 2: Schematic diagram of shoulder joint, 1: Fixed base, 2: Frame connecting rod, 3: Connecting rod, 4: Mobile platform

RRR linkages. Nine revolute axes of the mechanism intersect at one point O which is the center of rotation of the spherical body. The axes of three revolute joints connected with the fixed platform are vertical each other and the axes of three revolute joints connected with the moving platform are vertical to each other and the revolute joint axes of framed links and connecting rods are vertical to each other.

Kinematics analysis of the shoulder: Taking the center point O as the origin of the system coordinates, as shown in Fig. 2, a fixed Cartesian frame O{X, Y, Z} connected with the fixed platform and a moving Cartesian frame O{x, y, z} connected with the moving platform are established, with the X, Y, Z axes in coincidence with U_1, U_2, U_3 axes and the x, y, z axes in coincidence with V_1, V_2, V_3 axes, respectively. In the initial posture, the fixed Cartesian frame is coincident with the moving Cartesian frame.

The moving platform only can rotate round the fixed platform, referring to Zhen *et al.* (2006) the transformation matrix between the motion coordinate and the fixed coordinate is given as follows:

$$R = R_{ZYX}(\alpha, \beta, \gamma) = R(Z, \alpha)R(Y, \beta)R(X, \gamma) = \begin{bmatrix} \cos \alpha \cos \beta & \cos \alpha \sin \beta \sin \gamma - \sin \alpha \cos \gamma & \cos \alpha \sin \beta \cos \gamma + \sin \alpha \sin \gamma \\ \sin \alpha \cos \beta & \sin \alpha \sin \beta \sin \gamma + \cos \alpha \cos \gamma & \sin \alpha \sin \beta \cos \gamma - \cos \alpha \sin \gamma \\ -\sin \beta & \cos \beta \sin \gamma & \cos \beta \cos \gamma \end{bmatrix} \quad (1)$$

where, α, β and γ are the posture angles of the mobile platform.

When the moving platform pose changes, the direction cosine U_i, w_i and v_i of revolute joint of the i -th branch ($i = 1, 2, 3$) can be derived by:

$$\begin{cases} U_i = U_{i0} \\ V_i = R V_{i0} \\ W_i = |U_i \times V_i|^{-1} (U_i \times V_i) \end{cases} \quad (2)$$

where, U_{i0}, W_i and V_{i0} are the direction cosines of revolute joint of the i -th branch ($i = 1, 2, 3$) when the mechanism is on the initial state.

Referring to Zhen *et al.* (2006) study the constraint equation for the mechanism can be written as follows:

$$W_i \times V_t = 0 \quad (3)$$

Differentiating the constraint Eq. 3 with respect to time, yields:

$$\psi = J \dot{\theta} \quad (4)$$

where, $\dot{\theta} = (\dot{\theta}_1 \ \dot{\theta}_2 \ \dot{\theta}_3)$ is the vector of actuated joint rates and $\psi = (\dot{\alpha} \ \dot{\beta} \ \dot{\gamma})$ is the vector of the output angular velocities of the mechanism. J is the Jacobian matrix of the shoulder mechanism.

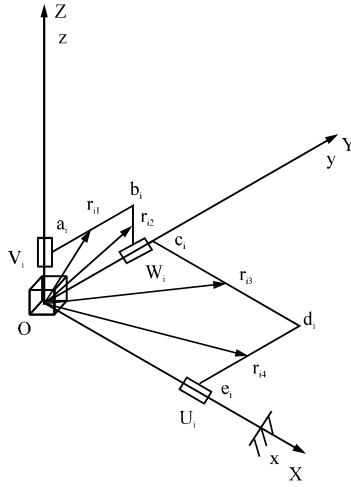


Fig. 3: Sketch of the i -th limb

Dynamic analysis: The schematic diagram of the i -th limb is shown in Fig. 3. The length of $a_i b_i$ and $d_i e_i$ is, l_1 , respectively, the length of $b_i c_i$ is l_2 , the length of $c_i d_i$ is l_3 , their linear density is ρ . Then the mass of $a_i c_i$ is $m_{ac} = \rho(l_1 + l_2)$, the mass of $c_i e_i$ is $m_{ce} = \rho(l_1 + l_3)$.

Kinetic energy analysis: The r_{11} , r_{12} , r_{14} and $d_i e_i$ are arbitrary dots in $a_i b_i$, $d_i e_i$, $d_i e_i$ and $b_i c_i$, respectively, as shown in Fig. 2 and their position vectors can be described as follows:

$$\begin{cases} r_{11} = l_2 V_i + x_1 W_i, & 0 \leq x_1 \leq l_1 \\ r_{12} = l_1 W_i + x_2 V_i, & 0 \leq x_2 \leq l_2 \\ r_{13} = l_1 W_i + x_3 U_i, & 0 \leq x_3 \leq l_3 \\ r_{14} = l_3 U_i + x_4 W_i, & 0 \leq x_4 \leq l_1 \end{cases} \quad (5)$$

The kinetic energies of $a_i b_i$, $b_i c_i$, $c_i d_i$ and $d_i e_i$ are:

$$\begin{cases} T_{iab} = \frac{1}{2} \rho \int_0^{l_1} \dot{r}_{11}^T \dot{r}_{11} dx_1 \\ T_{ibc} = \frac{1}{2} \rho \int_0^{l_2} \dot{r}_{12}^T \dot{r}_{12} dx_2 \\ T_{icd} = \frac{1}{2} \rho l_3 l_1^2 (\dot{W}_i^T \dot{W}_i) \\ T_{ide} = \frac{1}{6} \rho l_1^3 (\dot{W}_i^T \dot{W}_i) \end{cases} \quad (6)$$

The kinetic energy of the i -th limb is:

$$\begin{aligned} T_i &= T_{iab} + T_{ibc} + T_{icd} + T_{ide} \\ &= k_1 \dot{V}_i^T \dot{V}_i + k_2 \dot{W}_i^T \dot{W}_i + k_3 \dot{V}_i^T \dot{W}_i \end{aligned} \quad (7)$$

Where:

$$\begin{cases} k_1 = \frac{1}{2}\rho \left(l_1 l_2^2 + \frac{1}{3} l_2^3 \right) \\ k_2 = \frac{1}{2}\rho \left(\frac{2}{3} l_1^3 + l_1^2 l_2 + l_1^2 l_3 \right) \\ k_3 = \frac{1}{2}\rho (l_1^2 l_2 + l_1 l_2^2) \end{cases}$$

The total kinetic energy of the three limbs is:

$$T_L = \sum_{i=1}^3 T_i = \sum_{i=1}^3 (k_1 \dot{V}_i^T \dot{V}_i + k_2 \dot{W}_i^T \dot{W}_i + k_3 \dot{V}_i^T \dot{W}_i) \quad (8)$$

The moving platform only can rotate round the fixed platform, so its kinetic energy is:

$$T_D = \frac{1}{2} \dot{\Psi}^T I \dot{\Psi} \quad (9)$$

where, I is a inertia tensor of the moving platform. Based on theoretical mechanics, the following equation can be derived:

$$I = \frac{1}{6} \rho_0 l_2^5 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (10)$$

By substituting Eq. 10 into Eq. 9, the following equation is obtained:

$$T_D = k_4 \dot{\psi}^T \dot{\psi} \quad (11)$$

Where:

$$k_4 = \frac{1}{12} \rho_0 l_2^5$$

In view of Eq. 8 and 11, the total kinetic energy of the mechanism is:

$$T = T_L + T_D = \sum_{i=1}^3 (k_1 \dot{V}_i^T \dot{V}_i + k_2 \dot{W}_i^T \dot{W}_i + k_3 \dot{V}_i^T \dot{W}_i) + k_4 \dot{\psi}^T \dot{\psi} \quad (12)$$

Potential energy analysis: Let, position vectors of a_{c_1} and c_{e_1} be p_{i1} and p_{i2} , respectively, in view of Fig. 2, the position vectors are:

$$\begin{cases} p_{i1} = \begin{bmatrix} W_i & V_i & \frac{W_i \times V_i}{|W_i \times V_i|} \end{bmatrix} \begin{bmatrix} x_{ac} \\ y_{ac} \\ 0 \end{bmatrix} = x_{ac} W_i + y_{ac} V_i \\ p_{i2} = \begin{bmatrix} W_i & U_i & \frac{W_i \times U_i}{|W_i \times U_i|} \end{bmatrix} \begin{bmatrix} x_{ce} \\ y_{ce} \\ 0 \end{bmatrix} = x_{ce} W_i + y_{ce} U_i \end{cases} \quad (13)$$

where, $x_{ac} y_{ac}$ is the mass center coordinate of $a_i c_i$ in the $W_i O V_i$ plane, $x_{ce} y_{ce}$ is the mass center coordinate of $c_i e_i$ in the $W_i O U_i$ plane.

The potential energy of the i -th limb is:

$$E_i = m_{ac} g p_{i1z} + m_{ce} g p_{i2z} \quad (14)$$

Where:

$$\begin{aligned} p_{i1z} &= x_{ac} W_{iz} + y_{ac} V_{iz} \\ p_{i2z} &= x_{ce} W_{iz} + y_{ce} U_{iz} \end{aligned}$$

The total potential energy of the three limbs is:

$$\begin{aligned} E_1 &= \sum_{i=1}^3 E_i = m_{ac} g \left(\sum_{i=1}^3 p_{i1z} \right) + m_{ce} g \left(\sum_{i=1}^3 p_{i2z} \right) \\ &= k_5 (-\sin \alpha \sin \beta \sin \gamma + \cos \alpha \sin \gamma + \cos \alpha \cos \beta) + m_{ce} g y_{ce} \\ &\quad + m_{ac} g y_{ac} (\cos \beta \cos \gamma - \sin \beta + \cos \beta \sin \gamma) \end{aligned} \quad (15)$$

Where:

$$k_5 = m_{ac} g x_{ac} + m_{ce} g x_{ce}$$

Because the center of mass of the moving platform locates at the origin of the fixed coordinate system, the value of potential energy is 0. Then the total potential energy of the mechanism is:

$$E = E_1 \quad (16)$$

Dynamic modeling: Based on Lagrange method, the dynamic model of shoulder joint can be derived that:

$$f = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\psi}} \right) - \frac{\partial L}{\partial \psi} \quad (17)$$

where, $L = T - E$, in which T is kinetic energy, E is potential energy, $\psi = (\alpha \ \beta \ \gamma)^T$ is generalized coordinates, $f = (f_1 \ f_2 \ f_3)^T$ is generalized force.

By substituting Eq. 5-16 into Eq. 17, the following equation is obtained:

$$f = \frac{d}{dt} \left(\frac{\partial T_L}{\partial \dot{\psi}} \right) - \frac{\partial T_L}{\partial \psi} + \frac{d}{dt} \left(\frac{\partial T_D}{\partial \dot{\psi}} \right) - \frac{\partial T_D}{\partial \psi} - \frac{d}{dt} \left(\frac{\partial E}{\partial \dot{\psi}} \right) + \frac{\partial E}{\partial \psi} \quad (18)$$

The kinetic equation for the generalized coordinate α can be expressed as:

$$f_1 = k_1 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \dot{\alpha}} \right) - \frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \alpha} \right) + k_2 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \dot{\alpha}} \right) - \frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \alpha} \right) + k_3 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \dot{\alpha}} \right) - \frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \alpha} \right) + k_4 \frac{d}{dt} \left(\frac{\partial(\dot{\psi}^T \dot{\psi})}{\partial \dot{\alpha}} \right) + \frac{\partial E}{\partial \alpha} \quad (19)$$

The kinetic equation for the generalized coordinate β can be expressed as:

$$f_2 = k_1 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \dot{\beta}} \right) - \frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \beta} \right) + k_2 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \dot{\beta}} \right) - \frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \beta} \right) + k_3 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \dot{\beta}} \right) - \frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \beta} \right) + k_4 \frac{d}{dt} \left(\frac{\partial(\dot{\psi}^T \dot{\psi})}{\partial \dot{\beta}} \right) + \frac{\partial E}{\partial \beta} \quad (20)$$

The kinetic equation for the generalized coordinate γ can be expressed as:

$$f_3 = k_1 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \dot{\gamma}} \right) - \frac{\partial(\dot{V}_i^T \dot{V}_i)}{\partial \gamma} \right) + k_2 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \dot{\gamma}} \right) - \frac{\partial(\dot{W}_i^T \dot{W}_i)}{\partial \gamma} \right) + k_3 \sum_{i=1}^3 \left(\frac{d}{dt} \left(\frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \dot{\gamma}} \right) - \frac{\partial(\dot{V}_i^T \dot{W}_i)}{\partial \gamma} \right) + k_4 \frac{d}{dt} \left(\frac{\partial(\dot{\psi}^T \dot{\psi})}{\partial \dot{\gamma}} \right) + \frac{\partial E}{\partial \gamma} \quad (21)$$

Equation 18-20 can be rewritten into the general structure dynamic equation as:

$$f = D\ddot{\psi} + H[\dot{\psi}^2] + C[\dot{\psi} \dot{\psi}] + M \quad (22)$$

where, $\ddot{\psi} = (\ddot{\alpha} \ \ddot{\beta} \ \ddot{\gamma})^T$, $[\dot{\psi}^2] = [\dot{\alpha}^2 \ \dot{\beta}^2 \ \dot{\gamma}^2]^T$, $[\dot{\psi} \dot{\psi}] = [\dot{\alpha}\dot{\beta} \ \dot{\alpha}\dot{\gamma} \ \dot{\beta}\dot{\gamma}]^T$, D is the 3×3 symmetric positive equivalent inertia matrix, H is the 3×3 coefficient matrix of centrifugal force, C is the 3×3 coefficient matrix of coriolis force, M is the 3×1 gravity term.

Let input generalized force vector for parallel mechanism be $\tau = (\tau_1 \ \tau_2 \ \tau_3)^T$. The following equation is obtained via virtual work principle:

$$f = G\tau \quad (23)$$

where, G is force Jacobian matrix of the mechanism. According to dual relationship between motion and force transmissions in mechanism (Craig, 2011) yields:

$$G = J^{-1} \quad (24)$$

Substituting Eq. 24 and 23 into Eq. 22, the following equation is obtained, which is the inverse dynamics equation of the mechanism:

$$\tau = J(D\ddot{\psi} + H[\dot{\psi}^2]) + C[\dot{\psi} \dot{\psi}] + M \quad (25)$$

Equivalent inertia analysis: Equivalent inertia has effect on the stability, precision and dynamic response of control system which plays a leading role when the mechanism is accelerated and decelerated. The diagonal elements of the equivalent inertia matrix are also called effective inertia which can be used to estimate the equivalent inertia load of the servo motor. The matrix off-diagonal elements are also called coupling inertia, directly affect the stability of the control system. So, the equivalent inertia determines the dynamic characteristics of mechanism. It is necessary to analyze the equivalent inertia of the mechanism.

The equivalent inertia matrix of parallel shoulder can be derived from Eq. 25, as follows:

$$D = \begin{bmatrix} D_{11} & D_{12} & D_{13} \\ D_{21} & D_{22} & D_{23} \\ D_{31} & D_{32} & D_{33} \end{bmatrix} \quad (26)$$

Where:

$$D_{11} = 4k_1 + k_2(2\cos^2\alpha\sin^2\beta\cos^2\gamma + \sin 2\alpha\sin\beta\sin 2\gamma + 2\sin^2\alpha\sin^2\gamma + 2\sin^2\alpha\cos^2\beta + 2\sin^2\beta\sin^2\gamma + 2\cos^2\gamma) + 2k_4$$

$$D_{12} = D_{21} = k_2\left(\frac{1}{2}\sin 2\alpha\sin 2\beta\cos^2\gamma + \frac{1}{2}\sin 2\alpha\sin 2\beta - \cos^2\alpha\cos\beta\sin 2\gamma\right)$$

$$D_{13} = D_{31} = -4k_1\sin\beta - k_2\left(\frac{1}{2}\sin 2\alpha\sin^2\beta\sin 2\gamma + 2\cos^2\alpha\sin\beta\cos^2\gamma + 2\sin^2\alpha\sin\beta\sin^2\gamma\right) + \frac{1}{2}\sin 2\alpha\sin 2\gamma + 2\sin\beta$$

$$D_{22} = k_2(2\sin^2\beta\cos^2\gamma + 2\cos^2\alpha\cos^2\beta + 2\cos^2\beta\sin^2\gamma + 2\cos^2\beta + 2\sin^2\alpha\cos^2\beta\cos^2\gamma) + 2k_4 + 4k_1$$

$$D_{23} = D_{32} = k_2\left(\frac{1}{2}\sin 2\beta\sin 2\gamma + \frac{1}{2}\cos^2\alpha\sin 2\beta\sin 2\gamma - \sin 2\alpha\cos\beta\cos^2\gamma\right)$$

$$D_{33} = 4k_1 + k_2(2\cos^2\beta\sin^2\gamma + 2\cos^2\alpha\cos^2\gamma + 2\sin^2\alpha\sin^2\beta\sin^2\gamma + 2\sin^2\beta\cos^2\gamma + \sin 2\alpha\sin 2\gamma\sin\beta + 2\sin^2\gamma) + 2k_4$$

According to the discussion above, the equivalent inertia matrix which closely correlates with the sizes of mechanism and density of component is a time-varying function of mechanism configuration and changes with the body posture. This determines the dynamic characteristics of the mechanism.

RESULTS AND DISCUSSION

The structural parameters and inertial parameters of the shoulder are listed in Table 1. Let motion laws of moving platform of the shoulder joint be as follows:

$$\begin{cases} \alpha = \frac{\pi}{6} \sin(6t) & (\text{rad}) \\ \beta = \frac{\pi}{6} \cos(6t) & (\text{rad}) \\ \gamma = \frac{\pi}{6} \sin(6t) & (\text{rad}) \end{cases} \quad (27)$$

MATLAB (Matrix laboratory) is an interactive software system for numerical computations and graphics. It's a very powerful tool for doing numerical computations with matrices and vectors. Also, it has a variety of capabilities to display information graphically and can be extended through programs written in its own programming language. MATLAB provides a range of numerical computation methods for analyzing data, developing algorithms and creating models. So, the dynamic characteristics indices of the shoulder in task space can be obtained by a developed MATLAB program.

Solving inverse kinematics models, the curves for angular velocities of driving bars are obtained, as shown in Fig. 4. The angular velocities of driving bars regularly change with the

Table 1: Structural parameters and inertial parameters of the shoulder

Parameters	Value
Length of a_i, b_i and $d_i, e_{i11} = l_{14}$ (mm)	80
Length of b_i, c_{i112} (mm)	50
Length of c_i, d_i, l_{13} (mm)	150
Rod linear density ρ (kg m^{-1})	0.63
Moving platform density ρ (kg m^{-3})	2800
Acceleration of gravity ($\text{g m}^{-1} \text{sec}^{-2}$)	9.8

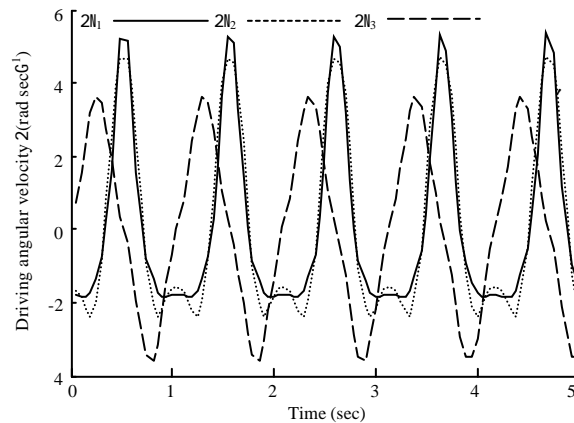


Fig. 4: Curves for driving angular velocities

posture parameters. The angular velocity maximum of the drive rod of branch 1 is $5.371 \text{ rad sec}^{-1}$, the angular velocity maximum of the drive rod of branch 2 is $3.628 \text{ rad sec}^{-1}$ and the angular velocity maximum of the drive rod of branch 3 is $4.689 \text{ rad sec}^{-1}$.

Solving the inverse dynamics Eq. 25, the curves for driving torques of driving bars are obtained, as shown in Fig. 5. In the process of movement, the driving torques of the drive rods regularly vary with the changing of posture parameters. The driving torque maximum of the drive rod of branch 1 is 0.4723 Nm , the driving torque maximum of the driving rod of branch 2 is 0.2281 Nm and the driving torque maximum of the drive rod of branch 3 is 0.3558 Nm .

Curves for inertia terms of the mechanism are shown in Fig. 6. It can be observed from Fig. 5 that the equivalent inertia parameters are functions of the parameters of the orientation and change with the mechanism posture. The trends of the equivalent inertia parameters and driving torques are consistent. The variation of the diagonal elements D_{11} and D_{33} of the inertia terms are same, their maximum and minimum values are same, respectively. The maximum value is

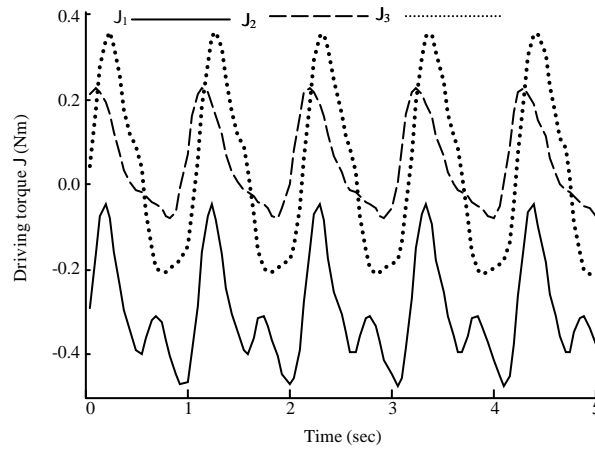


Fig. 5: Curves for driving torques

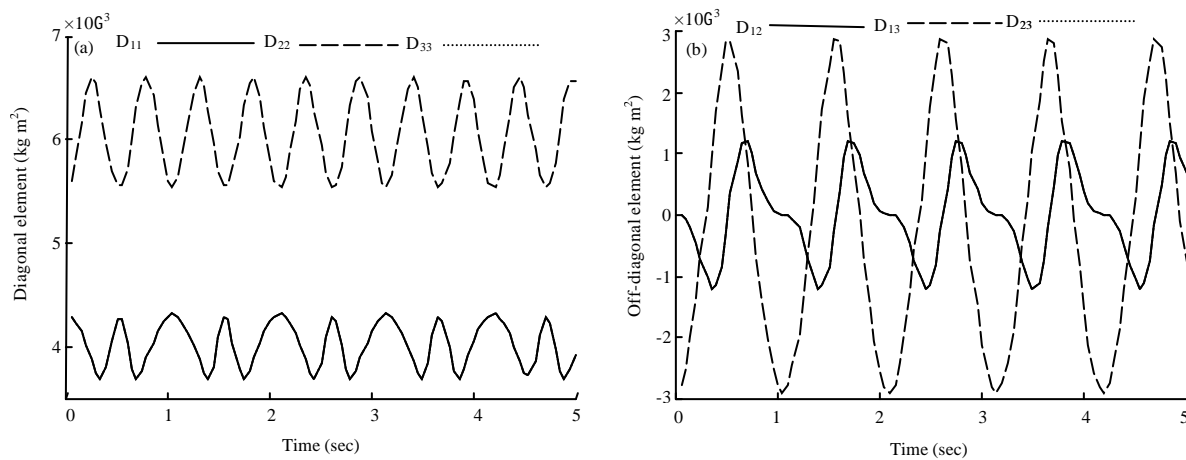


Fig. 6(a-b): Curves for inertia terms, (a) Diagonal element and (b) Off diagonal element

0.0043 kgm², the minimum value is 0.0036 kgm². The maximum value and minimum value of D22 is 0.0066 kgm² and 0.0055 kgm², respectively. The effective inertia of branch 2 is larger. The variation of the off-diagonal elements D12 and D23 of the inertia terms are same, their maximum and minimum values are same respectively, their maximum value is 0.0012 kgm², the maximum value of D13 is 0.0029 kgm². The coupling effect between the branch 1 and branch 3 is larger.

CONCLUSION

Through the analysis and calculation of kinetic energy and potential energy of each component of the novel 3-DOF robot shoulder joint, the dynamics model of the relationship between driving torque and the motion of the system parameters is established based on Lagrange method:

- On the basis of the dynamics model, the effective inertia, coupling inertia and driving torque of the parallel mechanism system are analyzed and the variation of the dynamics parameters of the mechanism is given, the dynamic coupling relationship between branches is analyzed with a developed MATLAB program
- The results show that, the angular velocity and driving torque maximum of the drive rod of branch 1 are the largest, the effective inertia of branch 2 is larger. The coupling effect between the branch 1 and branch 3 is larger and the posture change of the mechanism has a great influence on the dynamics parameters. The results can provide theoretical basis for the mechanism of dynamics optimization design, trajectory planning and control scheme

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